Thermodynamic properties of Nonadiabatic systems

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- Quantum Thermodynamics
 - Current Methods

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 - Define nonadiabatic

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 - measures the maximum or reversible work that may be performed by a thermodynamic system at a constant temperature and pressure



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- Equilibrium constant $K_{eq} = e^{\frac{-\Delta G}{RT}}$
- Molecular-scale systems $RT = k_B T \times N_A$



Quantum Thermodynamics

Partition function: $Z = \sum_{s} e^{-\beta E_s}$

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Thermodynamic properties expressed in terms of Z

$$A = -k_B T \ln(Z) \qquad C_v = \frac{1}{k_B T^2} \left(\frac{1}{Z} \frac{\partial^2 Z}{\partial \beta^2} - \left[\frac{-1}{Z} \frac{\partial Z}{\partial \beta} \right]^2 \right)$$
$$S = \frac{\partial}{\partial T} \left(k_B T \ln(Z) \right) \qquad U = -\frac{1}{Z} \frac{\partial Z}{\partial \beta} = -\frac{\partial \ln(Z)}{\partial \beta}$$



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 $Z = Z_t \times Z_R \times Z_N \times Z_v \times Z_e$

Z U, C_v, S, A, G

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Common approximations

1. Born Oppenheimer



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- 2. Harmonic oscillator



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- 3. Rigid rotor



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- 2. Harmonic oscillator
- 3. Rigid rotor
- 4. Ideal gas



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 - Options: ground state, DFT, B3LYP, cc-pVTZ

http://www.gaussian.com/g_whitepap/thermo.htm

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Ochterski, Joseph W. "Thermochemistry in Gaussian." *Http://www.gaussian.com*. Gaussian, Inc., 19 Apr. 2000. Web. http://www.gaussian.com/g_whitepap/thermo.htm.

Temperature	Pressure	
298.15 K	1 Atm.	

	Z	ln(Z)	E (KCal/Mol)	C _v (Cal/Mol-Kelvin)	S (Cal/Mol-Kelvin)
Electronic	1.000	0.000	0.000	0.000	0.000
Translational	0.583E+07	15.578	0.889	2.981	35.924
Rotational	51.266	3.937	0.592	1.987	9.811
Vibrational	1.000	0.007E-04	3.502	0.002	0.000
Total	0.299E+09	19.514	4.983	4.970	45.735
NIST source					45.7957 ± 0.001

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Limitations

• Systems in solution

Pes of 4 conformers of butane. Digital image. Http://chemistry.stackexchange.com/. N.p., 3 Oct. 2015. Web. 6 Apr. 2016. < http://i.stack.imgur.com/fgsHn.png>.

Limitations

- Systems in solution
- Molecules with many rotational conformers
 - Butane as shown below



Pes of 4 conformers of butane. Digital image. *Http://chemistry.stackexchange.com/*. N.p., 3 Oct. 2015. Web. 6 Apr. 2016. <http://i.stack.imgur.com/fgsHn.png>.

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Nonadiabatic systems

Systems of interest have many low lying electronic surfaces

transition metal atoms, radicals



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Common approximations:

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Systems of interest have many low lying electronic surfaces

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Common approximations:

Born Oppenheimer

Vibronic

- Harmonic oscillator
- **Rigid rotor** 3.

Ideal gas 4.



Born–Oppenheimer

$\hat{H} = \hat{T} + \hat{V}$
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$$\hat{H} = \hat{T}_N + \hat{T}_e + \hat{V}_{ee} + \hat{V}_{NN} + \hat{V}_{Ne}$$

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Electronic Hamiltonian

Nuclear & Electronic separation

 $\psi_{bo} = \phi(R) \chi(r;R)$ Nuclear Electonic

Nuclear & Electronic separation

$$\psi_{bo} = \underbrace{\phi(R)}_{\text{Nuclear Electonic}} \underbrace{\chi(r;R)}_{\text{Electonic}}$$

$$\hat{H}\psi_{bo} = \left[\hat{T}_N(R) + \hat{V}_{NN}(R) + \hat{H}_e(r;R)\right]\phi(R)\chi(r;R)$$

Obtain electronic energies

 $\hat{H}_e(r;R)\chi_n(r;R)$

$$= \left(\hat{T}_e(r) + \hat{V}_{ee}(r) + \hat{V}_{Ne}(r;R)\right) \chi_n(r;R)$$

 $= E_n(R)\chi_n(r;R)$

Obtain electronic energies

$\hat{H}_e(r;R)\chi_n(r;R)$



Solve for a given configuration R

$$\hat{H}\phi_n = \left[\hat{T}_N(R) + \underbrace{\hat{V}_{NN}(R) + E_e(R)}_{= \left[\hat{T}_N(R) + U_e(R)\right]\phi_n(R)} \phi_n(R)$$
$$= E\phi_n$$

Terminology

$$\psi_{bo} = \underbrace{\phi(R)}_{\text{Nuclear Electoric}} \underbrace{\chi(r;R)}_{\text{Electoric}}$$

Nuclear Electonic

• Adiabatic
$$\hat{H} = \begin{pmatrix} \hat{T}_{11} & \hat{T}_{12} \\ \hat{T}_{21} & \hat{T}_{22} \end{pmatrix} + \begin{pmatrix} \hat{V}_{11} & 0 \\ 0 & \hat{V}_{22} \end{pmatrix} \qquad \chi_A(r; R)$$

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$$\begin{array}{cccc} \bullet \mbox{ Adiabatic } \hat{H} = \begin{pmatrix} \hat{T}_{11} & \hat{T}_{12} \\ \hat{T}_{21} & \hat{T}_{22} \end{pmatrix} + \begin{pmatrix} \hat{V}_{11} & 0 \\ 0 & \hat{V}_{22} \end{pmatrix} & & \chi_A(r;R) \\ & & & & \\ \bullet \mbox{ Diabatic } \hat{H} = \begin{pmatrix} \hat{T}_{11} & \hat{T}_{12} \\ \hat{T}_{21} & \hat{T}_{22} \end{pmatrix} + \begin{pmatrix} \hat{V}_{11} & 0 \\ 0 & \hat{V}_{22} \end{pmatrix} & & \chi_D(r;R) \end{array}$$

Nonadiabatic systems

Unitary

R)

Transformation

Terminology

$$\psi_{bo} = \underbrace{\phi(R)}_{\text{Nuclear Electonic}} \underbrace{\chi(r;R)}_{\text{Electonic}}$$

• Adiabatic
$$\hat{H} = \begin{pmatrix} \hat{T}_{11} & \hat{T}_{12} \\ \hat{T}_{21} & \hat{T}_{22} \end{pmatrix} + \begin{pmatrix} \hat{V}_{11} & 0 \\ 0 & \hat{V}_{22} \end{pmatrix}$$

• Diabatic $\hat{H} = \begin{pmatrix} \hat{T}_{11} & 0 \\ 0 & \hat{T}_{22} \end{pmatrix} + \begin{pmatrix} \hat{V}'_{11} & \hat{V}_{12} \\ \hat{V}_{21} & \hat{V}'_{22} \end{pmatrix}$
• $\chi_D(r; R)$

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• $\chi_{D}(r; R)$

- Nonadiabatic
 - Systems with close or intersecting energy surfaces
 - Effects due to these off diagonal coupling terms

 Drives essential parts of visible and ultraviolet photochemistry and photobiology

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• Collisions of electronically excited species

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- Collisions of electronically excited species
- Chemiluminescent reactions
- Electron transfer processes

Conical intersections and internal conversion

<u>Internal conversion</u> Transition from higher to lower

electronic state without emission of photons



http://chem.libretexts.org/Core/Physical_and_Theoretical_Chemistry/Spectroscopy/Electronic_Spectroscopy/Jablonski_diagram

Sobolewski, Andrzej L., and Wolfgang Domcke. "The chemical physics of the photostability of life." Europhysics News 37.4 (2006): 20-23.

Conical intersections and internal conversion

<u>Conical intersection</u> Molecular Geometry point at which two potential energy surfaces are degenerate (intersect)

Suzuki, Toshinori. "Ultrafast Internal Conversion of Aromatic Molecules Studied by Photoelectron Spectroscopy using Sub-20 fs Laser Pulses." *Molecules* 19.2 (2014): 2410-2433.

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Conical intersection

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Ionization light

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Nonadiabatic effects

- Most discussion is in terms of dynamics and relates to processes that occur over time
- Our focus is on thermodynamic properties of nonadiabatic systems in thermal equilibrium



- Path Integral Formulation
 - Inclusion of electronic surfaces

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- Method 1: Direct calculation of Z
 - Importance Sampling
 - Monte Carlo Estimators

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- Method 1: Direct calculation of Z
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- Method 2: Coordinate Scaling

Feynman's Path Integrals

- An isomorphism between discrete quantum mechanical description and classical statistical mechanics of ring molecules
- Simulating classical systems are computationally advantageous

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Classical limit P = 1



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$$Z = \lim_{P \to \infty} \int d\{\mathbf{R}_{\gamma}\} \sum_{\{n_{\alpha}\}=1}^{N} \prod_{i=1}^{P} \langle \mathbf{R}_{i}, n_{i} | e^{-\tau \hat{H}} | \mathbf{R}_{i+1}, n_{i+1} \rangle$$

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Structure of Hamiltonian - Trotter

- Re-express vibronic Hamiltonian
- $\hat{H}_{vib} = \hat{T}(\mathbf{R}) + \hat{U}(\mathbf{R})$
- Trotter factorize

 $\langle \mathbf{R}_i, n_i | e^{-\tau \hat{H}} | \mathbf{R}_{i+1}, n_{i+1} \rangle$
• Re-express vibronic Hamiltonian

$$\hat{H}_{vib} = \hat{T}(\mathbf{R}) + \hat{U}(\mathbf{R}) \quad \blacksquare \quad \hat{H} = \hat{h}_o + \hat{V}$$

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$$\langle \mathbf{R}_{i}, n_{i} | e^{-\tau \hat{H}} | \mathbf{R}_{i+1}, n_{i+1} \rangle$$

$$\langle \mathbf{R}_{i}, n_{i} | e^{-\tau \hat{h}_{o}} | \mathbf{R}_{i+1}, n_{i} \rangle \langle \mathbf{R}_{i+1}, n_{i} | e^{-\tau \hat{V}} | \mathbf{R}_{i+1}, n_{i+1} \rangle$$

$$Harmonic oscillator propagator$$

$$\left[\mathbb{O}(\mathbf{R}_{i}, \mathbf{R}_{i+1}) \right]_{n_{i}, n_{i}}$$

$$\left[\mathbb{M}(\mathbf{R}_{i+1}) \right]_{n_{i}, n_{i+1}}$$

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• Trotter factorize

Analytical forms

• Probability distribution function of the system of interest $g(\{\mathbf{R}\}) = \operatorname{Tr}[\mathbb{O}(\mathbf{R}_1, \mathbf{R}_2)\mathbb{M}(\mathbf{R}_2)\mathbb{O}(\mathbf{R}_2, \mathbf{R}_3) \cdots \mathbb{O}(\mathbf{R}_P, \mathbf{R}_1)\mathbb{M}(\mathbf{R}_1)]$

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 $g({\mathbf{R}}) = \operatorname{Tr}[\mathbb{O}(\mathbf{R}_1, \mathbf{R}_2)\mathbb{M}(\mathbf{R}_2)\mathbb{O}(\mathbf{R}_2, \mathbf{R}_3)\cdots\mathbb{O}(\mathbf{R}_P, \mathbf{R}_1)\mathbb{M}(\mathbf{R}_1)]$

$$Z = \lim_{P \to \infty} \int \mathrm{d} \left\{ \mathbf{R}_{\gamma} \right\} g\left(\left\{ \mathbf{R} \right\} \right)$$

• Probability distribution function of a system with no vibronic coupling

$$\varrho(\{\mathbf{R}\}) = \mathrm{Tr}[\mathbb{O}(\mathbf{R}_1, \mathbf{R}_2)\mathbb{O}(\mathbf{R}_2, \mathbf{R}_3) \cdots \mathbb{O}(\mathbf{R}_P, \mathbf{R}_1)]$$

Expected value of function f(x) with normalized probability distribution function p(x) $\langle f(x) \rangle_p = \frac{1}{\int p(x)} \int p(x) f(x) \, \mathrm{d}x$

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$$\frac{1}{\int p(x) \, \mathrm{d}x} \int p(x) f(x) \, \mathrm{d}x = \frac{1}{N} \sum_{i=1}^{N} f(x_i) \pm \left(\frac{\sigma^2}{N}\right)^{\frac{1}{2}}$$

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Monte Carlo estimator

$$\langle f(x) \rangle_p = \lim_{N \to \infty} \frac{1}{N} \sum_{i=1}^N f(x_i)$$

Expected value of function f(x) with normalized probability distribution function p(x) $\langle f(x) \rangle_p = \frac{1}{\int p(x)} \int p(x) f(x) \, dx$



Importance Sampling

• Statistical method to reduce the variance of sampling

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Importance Sampling

- Statistical method to reduce the variance of sampling
- Sampling from a distribution p(x) is equivalent to sampling with weight $\frac{p(x)}{j(x)}$ from another distribution j(x)
- We would like to sample from $\varrho(\{\mathbf{R}\})$

$$Z = \int \mathrm{d} \left\{ \mathbf{R}_{\gamma} \right\} g\left(\left\{ \mathbf{R} \right\} \right)$$

$$Z = \int d\{\mathbf{R}_{\gamma}\} g(\{\mathbf{R}\})$$
$$Z = \int d\{\mathbf{R}_{\gamma}\} \varrho(\{\mathbf{R}\}) \frac{g(\{\mathbf{R}\})}{\varrho(\{\mathbf{R}\})} \frac{g(\{\mathbf{R}\})}{\varrho(\{\mathbf{R}\})}$$

$$Z = \int d\{\mathbf{R}_{\gamma}\}g(\{\mathbf{R}\})$$
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$$\frac{Z}{\int d\{\mathbf{R}_{\gamma}\}\rho(\{\mathbf{R}\})} = \frac{\int d\{\mathbf{R}_{\gamma}\}\rho(\{\mathbf{R}\})\frac{g(\{\mathbf{R}\})}{\rho(\{\mathbf{R}\})}}{\int d\{\mathbf{R}_{\gamma}\}\rho(\{\mathbf{R}\})}$$

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$$\frac{Z}{\int d\{\mathbf{R}_{\gamma}\} \varrho(\{\mathbf{R}\})} = \left| \left\langle \frac{g(\{\mathbf{R}\})}{\varrho(\{\mathbf{R}\})} \right\rangle_{\varrho} \right|$$

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$$Z = \left\langle \frac{g(\{\mathbf{R}\})}{\varrho(\{\mathbf{R}\})} \right\rangle_{\varrho} \left(\int d\{\mathbf{R}_{\gamma}\} \varrho(\{\mathbf{R}\})\right)$$

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$$A = -k_B T \ln(Z) \qquad C_v = \frac{1}{k_B T^2} \left(\frac{1}{Z} \frac{\partial^2 Z}{\partial \beta^2} - \left[\frac{-1}{Z} \frac{\partial Z}{\partial \beta} \right]^2 \right)$$
$$S = \frac{\partial}{\partial T} \left(k_B T \ln(Z) \right) \qquad U = -\frac{1}{Z} \frac{\partial Z}{\partial \beta} = -\frac{\partial \ln(Z)}{\partial \beta}$$

$$Z(\beta') = \int \mathrm{d}q \, \varrho(q,\beta) \frac{g(q',\beta')}{\varrho(q',\beta')}$$

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$$Cq = q' \quad C \,\mathrm{d}q = \mathrm{d}q'$$

$$Z(\beta') = \int dq \,\varrho(q,\beta) \frac{g(q',\beta')}{\varrho(q',\beta')} \qquad Cq = q' \quad C \, dq = dq'$$
$$\frac{\partial Z(\beta')}{\partial \beta'} = \int dq \,\varrho(q,\beta) \frac{\partial}{\partial \beta'} \left[f(q',\beta') \Big|_{\beta'=\beta} \right]$$
$$= \int dq \,\varrho(q,\beta) \frac{1}{2\Delta\beta} \left[f(q',\beta+\Delta\beta) - f(q',\beta-\Delta\beta) \right]$$

$$\begin{split} Z(\beta') &= \int \mathrm{d}q \, \varrho(q,\beta) \frac{g(q',\beta')}{\varrho(q',\beta')} \qquad Cq = q' \quad C \, \mathrm{d}q = \mathrm{d}q' \\ \frac{\partial Z(\beta')}{\partial \beta'} &= \int \mathrm{d}q \, \varrho(q,\beta) \frac{\partial}{\partial \beta'} \left[f(q',\beta') \right|_{\beta'=\beta} \\ &= \int \mathrm{d}q \, \varrho(q,\beta) \frac{1}{2\Delta\beta} \Big[f(q',\beta+\Delta\beta) - f(q',\beta-\Delta\beta) \Big] \\ \frac{\partial^2 Z}{\partial \beta^2} &= \int \mathrm{d}q \, \varrho(q,\beta) \frac{1}{\Delta \beta^2} \Big[f(q',\beta+\Delta\beta) + f(q',\beta-\Delta\beta) - 2f(q',\beta) \Big] \end{split}$$

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$$A = \langle e \rangle_g - T \left[\int_0^{T_f} \frac{\langle c \rangle_g}{T} \, \mathrm{d}T \right] \qquad C_v = \langle c \rangle_g = \frac{\langle cw \rangle_\varrho}{\langle w \rangle_\varrho}$$
$$S = \int_0^{T_f} \frac{C_v}{T} \, \mathrm{d}T \qquad \qquad U = \langle e \rangle_g = \frac{\langle ew \rangle_\varrho}{\langle w \rangle_\varrho}$$

Sum over states method

Armadillo C++ template library

$$Z = \sum_{n} e^{-\beta E_n}$$

$$A = U - TS$$

$$C_v = \frac{1}{k_B T^2} \left(\sum_n E_n^2 \frac{e^{-\beta E_n}}{Z} - U^2 \right)$$

$$U = \frac{1}{Z} \sum_n E_n e^{-\beta E_n} \quad S = -\frac{k_B}{Z} \sum_n e^{-\beta E_n} \ln(p_n)$$

Sum over states method

• Vibronic Hamiltonian (with coupling)

$$\begin{bmatrix} E_{1} + \frac{1}{2}\omega_{1}(q - q_{m1})^{2} & \mu q & \mu q & \mu q \\ \mu q & E_{2} + \frac{1}{2}\omega_{2}(q - q_{m2})^{2} & \mu q & \mu q \\ \mu q & \mu q & E_{3} + \frac{1}{2}\omega_{3}(q - q_{m3})^{2} & \mu q & \mu q \\ \mu q & \mu q & \mu q & E_{4} + \frac{1}{2}\omega_{4}(q - q_{m4})^{2} & \mu q \\ \mu q & \mu q & \mu q & \mu q & E_{5} + \frac{1}{2}\omega_{5}(|q - q_{m5})^{2} \end{bmatrix}$$

Sum over states method

• Vibronic Hamiltonian (with coupling)

 $\begin{bmatrix} E_{1} + \frac{1}{2}\omega_{1}(q - q_{m1})^{2} & \mu q & \mu q & \mu q \\ \mu q & E_{2} + \frac{1}{2}\omega_{2}(q - q_{m2})^{2} & \mu q & \mu q \\ \mu q & \mu q & E_{3} + \frac{1}{2}\omega_{3}(q - q_{m3})^{2} & \mu q & \mu q \\ \mu q & \mu q & \mu q & E_{4} + \frac{1}{2}\omega_{4}(q - q_{m4})^{2} & \mu q \\ \mu q & \mu q & \mu q & \mu q & E_{5} + \frac{1}{2}\omega_{5}(|q - q_{m5})^{2} \end{bmatrix}$

Harmonic Hamiltonian (without coupling)

$$\begin{bmatrix} E_1 + \frac{1}{2}\omega_1 (q - q_{m1})^2 & 0 & 0 & 0 \\ 0 & E_2 + \frac{1}{2}\omega_2 (q - q_{m2})^2 & 0 & 0 & 0 \\ 0 & 0 & E_3 + \frac{1}{2}\omega_3 (q - q_{m3})^2 & 0 & 0 \\ 0 & 0 & 0 & E_4 + \frac{1}{2}\omega_4 (q - q_{m4})^2 & 0 \\ 0 & 0 & 0 & 0 & E_5 + \frac{1}{2}\omega_5 (q - q_{m5})^2 \end{bmatrix}$$





• Completed PI python script for uncoupled system

$$Z = \left\langle \frac{g\left(\{\mathbf{R}\}\right)}{\varrho\left(\{\mathbf{R}\}\right)} \right\rangle_{\varrho} \left(\int \mathrm{d}\left\{\mathbf{R}_{\gamma}\right\} \varrho\left(\{\mathbf{R}\}\right) \right)$$
Current Progress

• Completed PI python script for uncoupled system

$$Z = \left\langle \frac{g\left(\{\mathbf{R}\}\right)}{\varrho\left(\{\mathbf{R}\}\right)} \right\rangle_{\varrho} \left(\int \mathrm{d}\left\{\mathbf{R}_{\gamma}\right\} \varrho\left(\{\mathbf{R}\}\right) \right)$$

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Current Progress

Completed PI python script for uncoupled system

$$Z = \left\langle \frac{g\left(\{\mathbf{R}\}\right)}{\varrho\left(\{\mathbf{R}\}\right)} \right\rangle_{\varrho} \left(\int \mathrm{d}\left\{\mathbf{R}_{\gamma}\right\} \varrho\left(\{\mathbf{R}\}\right) \right)$$

 $g({\mathbf{R}}) = \operatorname{Tr}[\mathbb{O}(\mathbf{R}_1, \mathbf{R}_2)\mathbb{M}(\mathbf{R}_2)\mathbb{O}(\mathbf{R}_2, \mathbf{R}_3)\cdots\mathbb{O}(\mathbf{R}_P, \mathbf{R}_1)\mathbb{M}(\mathbf{R}_1)]$

 $\varrho(\{\mathbf{R}\}) = \operatorname{Tr}[\mathbb{O}(\mathbf{R}_1, \mathbf{R}_2) \mathbb{O}(\mathbf{R}_2, \mathbf{R}_3) \cdots \mathbb{O}(\mathbf{R}_P, \mathbf{R}_1)]$

Current Progress



$\frac{g\left(\{\mathbf{R}\}\right)}{\varrho\left(\{\mathbf{R}\}\right)}\right\rangle_{\varrho}$

Future

- Implement support for arbitrary coupling propagator $[\mathbb{M}(\mathbf{R}_{i+1})]_{n_i,n_{i+1}}$
- Implement quadratic and linear terms in SOS/analytical code(*correctly)
- Implement analytic portion of Z calculation

$$\left(\int \mathrm{d}\left\{\mathbf{R}_{\gamma}\right\}\varrho\left(\left\{\mathbf{R}\right\}\right)\right)$$

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• Replicate the Python code in Fortran for increased performance

Conclusions

- Developed a path integral formulation for nonadiabatic systems
- Developed two approaches to calculating Z
- Derived estimators for thermodynamic properties
- About to start testing convergence of Path Integral Code